Short note

The γ -Decays of ²¹⁰Po-Levels from the (³He, d* $\gamma\gamma$)-Reaction

H. Klein¹, I. Wiedenhöver^{1,a}, H. Tiesler¹, H. Meise¹, A. Fitzler¹, A. Dewald¹, H.G. Thomas¹, D. Weißhaar¹, P. von Brentano¹

Institut für Kernphysik, Universität zu Köln, Zülpicher Strasse 77, D-50937 Köln

Received: 17 December 1998 / Revised version: 27 January 1999 Communicated by D. Schwalm

Abstract. An in-beam experiment with the subcoulomb reaction $^{209}\text{Bi}(^{3}\text{He}, d^{*}\gamma\gamma)^{210}\text{Po}$ at 20.5 MeV was performed with two EUROBALL CLUSTER detectors in Cologne. It closed the gap between the low energy levels of the level-scheme and the high energy levels found in $^{209}\text{Bi}(^{3}\text{He}, d)^{210}\text{Po}$ and $^{208}\text{Pb}(^{4}\text{He}, t)^{210}\text{Po}$ particle experiments. New branchings have been found and the $(^{3}\text{He}, d^{*}\gamma\gamma)$ reaction below the coulomb-barrier has been used successfully.

PACS. 21.10.Pc Single-particle levels and strength function -23.20.Lv Gamma transitions and level energies -25.55.-e 3H-,3He-, and 4He-induced reactions

The mass region around the doubly magic nucleus ²⁰⁸Pb is one of the most important testing grounds for the nuclear shell model [1]. Although the excitations of ²⁰⁹Bi(³He, d)²¹⁰Po and ²⁰⁸Pb(⁴He, t)²¹⁰Po have been well investigated with methods of particle spectroscopy [2], the electromagnetic decay properties of these states are not known to the same extent. No γ -transitions depopulating the multiplet-levels above 4 MeV had been observed so far [3]. The observation of these high energy γ -transitions and thus the connection to the level-scheme has become possible due to the high efficiency of the EUROBALL CLUSTER detectors.

The experiment, performed at the Cologne FN-Tandem accelerator, used the ²⁰⁹Bi(³He, d* $\gamma\gamma$)²¹⁰Po reaction at 20.5 MeV. We write d* to include the various exit channels of deuterons with spin 1 and 0 and the (pn)-channel, which we did not distinguish in particular. Thus we have d (S=1), d (S=0), (pn) \subset d*. In order to reduce the competing fusion reaction (³He, 2n) we chose a beam energy below the Coulomb-barrier. This reaction was also chosen to test the power of the subcoulomb stripping reaction in gamma decay studies. We were able to observe γ -spectra with only relatively few contaminations, which were mainly (³He, 2n)²⁰⁹Po (~50% compared with ²¹⁰Po), (³He, 3n)²⁰⁹At (~10%) and (³He, 2n)²¹⁰At (~5%). These contaminations could easily be separated using $\gamma\gamma$ -coincidences.

The spectrometer used in this experiment consisted of two unshielded EUROBALL CLUSTER detectors at a distance to the target of only 8 cm. Each of the EU-ROBALL CLUSTER detectors, developed by a collaboration of groups from Cologne, Jülich, the company EURISYS, and the EUROBALL Collaboration [4] consists of seven high purity germanium detectors (segments) with an efficiency of 65% relative to a 3"x3" NaI detector. The setup was optimized to detect $\gamma\gamma$ -coincidence events of high energy transitions. The close target-detector distance is possible, because our proton stripping reactions provide γ cascades of low multiplicities, which reduces the probability of multiple hits respective to fusion-evaporation reactions. With the two detectors we achieved a photopeak efficiency of \approx 7% at $1333 \,\mathrm{keV}$, applying the nearest neighbour addback method as described in [5]. A schematic view of the setup can be seen in Fig. 1.

The granularity of the EUROBALL CLUSTER detector allows a quantitative analysis of γ -peaks due to sum-up events [6]. To obtain a good energy and efficiency calibration at high energies, in addition to the standard ²²⁶Ra source we used ²⁴Mg (²⁴Mg(p,n)²⁴Al $\xrightarrow{\beta+}$ ²⁴Mg* $\xrightarrow{\gamma}$ ²⁴Mg) for calibration. This procedure [7] allows a highest quality energy calibration up to 7 MeV. Due to the different decay-multiplicities of ²²⁶Ra, ²⁴Mg and ²¹⁰Po a systematic error in the efficiency calibration had to be taken into account. This error emerged to be less than 10%.

The $\gamma\gamma$ -matrix of addback events contains 1.2 billion coincidence events from 40 hours of measuring. This data were analyzed to establish the coincidence relations and extend the levelscheme. We were able to find 14 new lev-

^{*a*} Present address: Argonne National Laboratory, 9700 South Cass Avenue, Argonne, Illinois 60439, USA



Fig. 1. Experimental setup: Two EUROBALL CLUSTER detectors in a distance of 8 cm from the target

Table 1. Comparison of sets of levels observed by different $^{210}\mathrm{Po}$ reactions

	reaction	Ref.	levels
1:	$^{209}_{208}\text{Bi}(^{3}\text{He}, d)$	[3]	23
2:	208 Po(⁴ He, t)	[3]	31
3:	209 Bi(t, 2n γ)	[8]	45
4:	$^{209}\text{Bi}(^{3}\text{He},\mathrm{d}^{*}\gamma\gamma)$	this work	56

els and 39 new γ -transitions in ²¹⁰Po. The results of our analysis are shown in Fig. 2 and in Table 2. Furthermore, six levels previously observed only in particle experiments were connected to the level scheme by γ -transitions and thus the precision of the measured level energies could be improved.

This experiment gives a good evidence for the utility of the subcoulomb (³He, d* $\gamma\gamma$) reaction. The amount of new data is remarkable taking into account the previous knowledge of 61 levels [3]. A comparison of different reactions used to investigate states of ²¹⁰Po is given in Tab. 1.

In Fig. 2 the single particle excitations in ²¹⁰Po are compared with those of ²⁰⁹Bi [2]. We see all states of these multiplets with exception of a (3⁺) state at 4320 MeV. Because of the observed γ -transitions, it is now possible to improve several spin assignments. The 4553 keV state was (4⁺,7⁺) and is now assigned (7⁺) due to the observed γ -transition to 8⁺ states. Thus it belongs to the $h_{9/2} \otimes p_{3/2}$ multiplet. The missing 4⁺ could not be found.

The 3^- state at 2846 keV was assigned to the $h_{9/2} \otimes i_{13/2}$ multiplet, because it is the only 3^- in this region. The spin and parity of the 3^+ at 4592 keV assigned by [2] are put in some doubt because here an M3 would compete with possible E1, M1 and E2 transitions, which were not observed.

The fact, that the (³He, pn)-reaction populated nearly all single-particle excitations may give a hint, that a subcoulomb stripping of one proton of the ³He occured and the output-channel was a deuteron d (S=1) or an excited deuteron d (S=0), which would decay to the (pn)-channel. Furthermore, we observed spin transfers up to $6\hbar$.



Fig. 2. New γ -transitions found in ²¹⁰Po. Only transitions which complete the 2-particles multiplets [2] are shown

Besides the construction of a level scheme with the new γ -transitions, the decay branchings of the levels were found. To obtain high-quality branching ratios from our $\gamma\gamma$ -matrices, we gate on transitions populating the level of interest. This was not possible for the new high energy levels. The branchings of these levels were found in the gate of the $2_1^+ \rightarrow 0_1^+$ 1181 keV transition. This includes additional systematic errors of about 30%.

The branchings will provide a useful test of the shellmodel. There is still a need for further experiments to obtain spin and parity assignments, since some of the known spins are only tentative and our experimental setup offered no possibility to analyse angular distribution or correlation.

Summing up, this experiment attests the spectroscopic power of the subcoulomb (³He, d* $\gamma\gamma$)-reaction in conjunction with two EUROBALL CLUSTER detectors. Thus 56 new levels of ²¹⁰Po were observed up to an energy of 5 MeV.

The authors thank Dr. J. Eberth, O. Thelen and A. Lisetzkiy for fruitful discussions. Furthermore, we thank Drs. A. Jungclaus and C. Teich from the University of Goettingen for the help in setting up a EUROBALL CLUSTER detector. We also thank the referee for constructive remarks. This work was supported by DFG under project no. Br 799/7-1.

Table 2. New γ -transitions and energy-levels observed in ²¹⁰Po. If published energy E_p of the levels are given, the transition energies E_{γ} and the branchings are new. Initial and final spins (J_i, J_f) and the final levels E_f were taken from Ref. [3]. Changes from these values are marked with an asterix *. The relative intensities were calculated from all γ -transitions depopulating a level, seen in coincidence with the $2_1^+ \rightarrow 0_1^+$ transition. The given relative cross-sections $I_{\rm rel}$ are taken from the (³He, d)-particle experiment [2]. The intensities and cross-sections are normalized to the 4141 keV level

$E_i \; [keV]$	$E_p \; [keV]$	\mathbf{J}_i^{π}	\mathbf{J}_{f}^{π}	$E_{\gamma} \ [keV]$	E_f [keV]	Branching	$\mathrm{I}_{\mathrm{rel}}$	$\sigma_{ m rel40^\circ}$
3477.3_{3}			8+	1289.3_{2}	2188.0	100	9	
3637.5_{2}			3 -	1250.7_{2}	2386.8	100	12	
3693.9_{2}			3 -	1307.1_{2}	2386.8	100	5	
—	3792_{2}	(2+)						2
—	4027_{1}	4 +						45
4029.1_{3}			8 +	2602.4_3	1557.0	100	30	
4043.4_{3}			8 +	1855.4_{2}	2188.0	100	25	
4105.1_{3}	4105_{10}		8+	1917.1_{2}	2188.0	100	5	
4141.1_{4}	4139_{1}	(6+)	8+	2665.5_{4}	1473.3	30 10	100	100
			7+	1702.5_{2}	2438.4	15_{-10}		
			8+	1953.6_{2}	2188.0	30 10		
			8 +	2583.8_{3}	1557.0	100 20		
			8+	2671.4_{9}	1473.3			
—	4237_{10}							
—	4320_{1}	(3+)						56
4329.5_{4}			6+	2003.5_{3}	2326.0	100	15	
—	4382_{1}	(5+)						88
4386.9_{4}			6+	2059.9_{5}	2326.0	40 15	55	
			8+	2199.3_{3}	2188.0	100 10		
4469.6_{3}	4469_{1}	(6+)	6+	2143.5_{3}	2326.0	10 s	120	170
			8+	2281.9_{3}	2188.0	20 10		
			8+	2913.1_{3}	1557.0	100 15		
4542.1_{3}		$(4+)^*$	5+	2139.2_{3}	2403.3	60 20	35	
			3 -	2159.8_{3}	2382.6	40 10		
			4 +	2215.7_{4}	2326.0	50 25		
			6+	3115.6_{6}	1426.7	100 зо		
4553.6_{4}	4553_{1}	$(7+)^*$	8+	2365.6_{4}	2188.0	80 20	70	210
			8+	2997.9_{6}	1557.0	100 20		
4592.5_{4}	4591	(3+)	0+	4592.5_{4}	0.0	100		630
4621.5_{3}		$(3+)^*$	5+	2207.9_{3}	2413.8	100 зо	85	
			1+	2227.7_{3}	2393.8	30 10		
			4 +	2238.8_{4}	2382.6	45 15		
			2 +	2331.5_{3}	2290.1	50 25		
_	4624_{1}	(5+)						
4637.5_{3}			5+	2234.7_{4}	2403.3	60 <i>зо</i>	90	
			4 +	2255.1_{3}	2382.6	100 15		
			6+	2311.5_{4}	2326.0	25 10		
4644.9_{5}	4644_{1}	(6+)	8 +	2456.9_{5}	2188.0	100	5	
4660.4_{3}			4 +	2277.8_{3}	2382.6	75 25	50	
			6+	2334.1_{4}	2326.0	100 40		
4948.13			5+	2544.8_{3}	2403.3	50 25	45	
			6+	3474.9_{5}	1473.3	100 50		
4974.4_{5}			8 +	2786.4_{5}	2188.0	100	10	
_	4991_{1}							80
4998.2_{5}			8+	2810.2_{5}	2188.0	100	2	

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